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## Numerical Investigations of the Impact of Peclet Number on the Transient Nonlinear Thermal Response of Radiating-Convecting Internally-Heated Porous Moving Fin using Method of Lines

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### ABSTRACT

In the present study, numerical method of lines is applied to study the impact of Peclet number on the transient nonlinear thermal behaviour of moving porous fins. Through energy analysis of the passive device, the transient thermal model of the fin is developed. The nonlinear partial differential equation is nondimensionalized to directly establish the Peclet number in the adimensional thermal model. Thereafter, the model is solved by numerical method of lines and the influences of Peclet number on the thermal response of the fins is explored. The results of the numerical investigations illustrate that the temperature of the fin is enhanced when the Peclet number is augmented. This shows that low value of Peclet number favours cooling enhancement. Also, under varying Peclet number, the extended surface thermal distribution decreases as porous, conductive-convective, conductive-radiative and porous terms increase. However, the temperature of the passive device rises as internal heat generation, ambient and surface temperatures is heightened. This study has given better physical insights and understanding of the thermal problems in extended surfaces.

**Keywords:** Moving fin, Peclet number, Porous Fin, Thermal analysis, Transient response, Numerical method of lines

## **1. INTRODUCTION**

The applications of fins and spines are for augmentations of the rate of heat transfer in thermal and electronic components. The theoretical investigations of fin behaviour from to the nonlinear thermal models have been analyzed in different works. In such studies, Torabi et al. [1] Aziz and Khani [2] presented analytical solution for continuously moving convecting-radiating fin with variable thermal conductivity using differential transformation and homotopy analysis methods while Aziz and Lopez [3] performed such thermal analysis for sheet/rod.

Singh et al. [4] applied wavelet collocation solution for the same problem. Aziz and Torabi [5] studied the passive device with simultaneous variation of thermal properties. Ma et al. [6] adopted spectral collocation method for radiating–conducting porous fin with variable thermal properties.

Sun et al. [7] utilized the same method for the thermal analysis of a moving rod with variable thermal conductivity.

Kanth and Kumar [8, 9] used Haar wavelet method for the moving fin with variable thermal conductivity. Adomian decomposition method was made use by Singla and Ranjan [10]. Moradi and Rafiee [11] as well as Dogonchi and Ganji [12] explored differential transformation method (DTM) to generate for the thermal behaviour of continuously moving convecting-radiating fin with variable thermal conductivity. Sun and Ma [13] implemented collocation spectral method for moving convecting-radiating fin with variable thermal properties and internal heat generation.

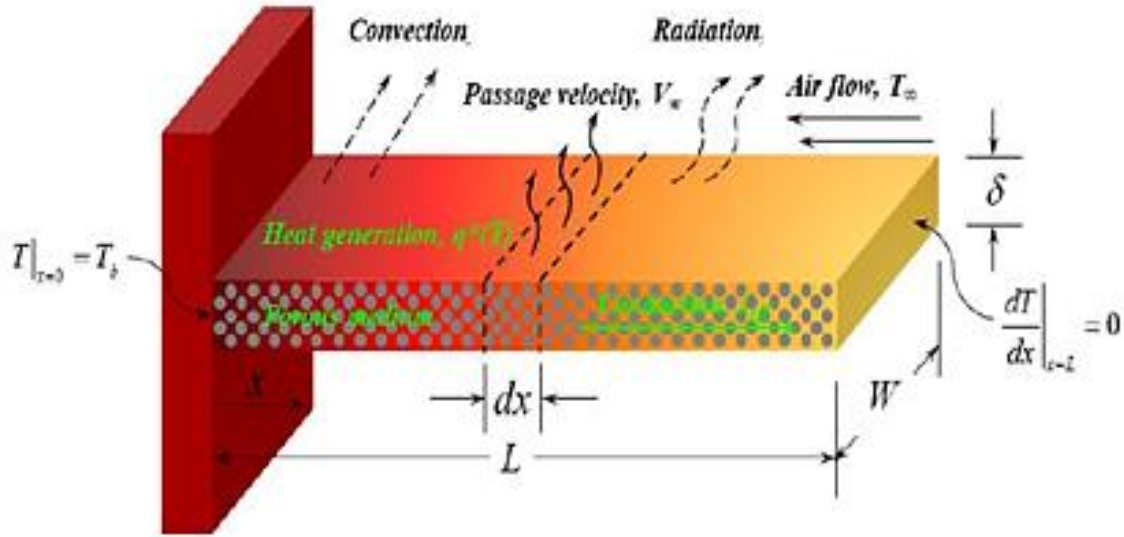
In their studies, Sobamowo et al. [14] compared differential transformation and finite difference methods for the thermal problems while Ndlovu and Moitsheki [15-17] used DTM to study transient thermal problem through a moving fin with variable thermal properties. Other researchers [18-37] adopted various analytical, numerical, semi-analytical, semi-numerical and approximate analytical methods to study heat transfer in continuously moving device. However, a study on the exploration of Peclet number effects on the on the thermal responses of fins is limited in literature.

Therefore, in this work, numerical method of lines is adopted to solve the transient thermal model of a convective-radiative rectangular moving porous fin with temperature dependent thermal conductivity and internal heat generation. The numerical solutions are used to established the effect of Peclet number on the thermal behaviour of the fin.

## **2. MODEL DEVELOPMENT FOR THE TRANSIENT THERMAL FLOW PROCESS**

Consider a porous moving fin which is internally heated. Given that the fin has length  $L$ , thickness  $\delta$  and perfectly and thermally attached to a prime surface at temperature  $T_b$ .

The various assumptions for the present study can be found in our previous works [14, 29].



**Fig. 1.** Schematic of a longitudinal moving porous fin with perfect thermal contact and insulated tip

Adopting the above model assumptions, the energy analysis from Fig. (1) provides the differential equation governing the thermal model as

$$\frac{\partial}{\partial x} \left( k_{eff} \frac{\partial T}{\partial x} \right) + \frac{4\sigma\phi}{3\beta_R} \frac{\partial}{\partial x} \left( \frac{\partial T^4}{\partial x} \right) - \frac{\rho_f c_{p,f} g K \beta \phi (T - T_a)^2}{v\delta} - \frac{hP(1-\phi)}{A_{cr}} (T - T_a) - \frac{\sigma \varepsilon P}{A_{cr}} (T^4 - T_s^4) + (1-\phi)q'''(T) = (\rho c_p)_{eff} u \frac{\partial T}{\partial x} + (\rho c_p)_{eff} \frac{\partial T}{\partial t} \quad .. (1)$$

The variable thermal conductivity and internal heat generation are respectively given by the linear expressions as

$$k_{eff} = (1-\phi)k_s + \phi k_f \quad (2)$$

where,

$$q'''(T) = q_o''' [1 + \gamma(T - T_a)] \quad (3)$$

and

$$(\rho c_p)_{eff} = (1-\phi)(\rho c_p)_s + \phi(\rho c_p)_f \quad (4)$$

Then, the governing equation becomes

$$k_{eff} \frac{\partial^2 T}{\partial x^2} + \frac{4\sigma\phi}{3\beta_R} \frac{\partial}{\partial x} \left( \frac{\partial T^4}{\partial x} \right) - \frac{\rho_f c_{p,f} g K \beta \phi (T - T_a)^2}{v\delta} - \frac{h_b P (1 - \phi) (T - T_a)}{A_{cr}} \tag{5}$$

$$- \frac{\sigma \epsilon_b P}{A_{cr}} (T^4 - T_s^4) + (1 - \phi) q_o''' [1 + \gamma(T - T_a)] = (\rho c_p)_{eff} \left[ u \frac{\partial T}{\partial x} + \frac{\partial T}{\partial t} \right]$$

Linearizing the radiative term based on Roseland’s approximation, we expanded  $T^4$  about  $T_a$  with the aid of Taylor series,

$$T^4 \cong T_a^4 + 4T_a^3 (T - T_a) + 6T_a^2 (T - T_a)^2 + 4T_a (T - T_a)^3 + \dots \tag{6}$$

If the higher order components in Eq. (6) are ignored, we have

$$T^4 \cong 4T_a^3 T - 3T_a^4 \tag{7}$$

Therefore,

$$T^4 - T_a^4 \cong 4T_a^3 (T - T_a) \tag{8}$$

Putting Eq. (8) into Eq. (5),

$$\frac{4\sigma\phi}{3\beta_R} \frac{\partial T^4}{\partial x} = \frac{4\sigma\phi}{3\beta_R} \frac{\partial (4T_a^3 T - 3T_a^4)}{\partial x} = \frac{16\sigma\phi T_a^3}{3\beta_R} \frac{\partial T}{\partial x} \tag{9}$$

Putting Eqs. (8) and (9) in Eq. (5),

$$k_{eff} \frac{\partial^2 T}{\partial x^2} + \frac{16\sigma\phi T_a^3}{3\beta_R} \frac{\partial^2 T}{\partial x^2} - \frac{\rho_f c_{p,f} g K \beta W \phi (T - T_a)^2}{vA_{cr}} - \frac{hP(1 - \phi)}{A_{cr}} (T - T_a) \tag{10}$$

$$- \frac{4\sigma\epsilon P T_a^3}{A_{cr}} (T - T_a) + (1 - \phi) q_o''' [1 + \gamma(T - T_a)] = (\rho c_p)_{eff} u \frac{\partial T}{\partial x} + (\rho c_p)_{eff} \frac{\partial T}{\partial t}$$

The initial condition is

$$T = T_0, \text{ when } t = 0, \text{ for } 0 < x < L, \tag{11}$$

The boundary conditions

$$T = T_b, \text{ at } x = 0, \text{ for } t > 0, \tag{12}$$

$$\frac{dT}{dx} = 0, \text{ at } x = L, \text{ for } t > 0, \tag{13}$$

Eq. (10) can be expressed as

$$\left(k_{eff} + \frac{16\sigma\phi T_a^3}{3\beta_R}\right) \frac{\partial^2 T}{\partial x^2} - \frac{\rho_f c_{p,f} g K \beta W \phi (T - T_a)^2}{v A_{cr}} - \frac{hP(1-\phi)}{A_{cr}} (T - T_a) - \frac{4\sigma\varepsilon P T_a^3}{A_{cr}} (T - T_a) + (1-\phi) q_o''' [1 + \gamma(T - T_a)] = (\rho c_p)_{eff} u_o \frac{\partial T}{\partial x} + (\rho c_p)_{eff} \frac{\partial T}{\partial t} \tag{14}$$

The variable internal heat generation is stated as follows

$$q'''(T) = (1-\phi) q_o''' + (1-\phi) q_o''' \gamma T - (1-\phi) q_o''' \gamma T_a \tag{15}$$

Therefore

$$\left(k_{eff} + \frac{16\sigma\phi T_a^3}{3\beta_R}\right) \frac{\partial^2 T}{\partial x^2} - \frac{\rho_f c_{p,f} g K \beta W \phi (T - T_a)^2}{v A_{cr}} - \frac{hP(1-\phi)}{A_{cr}} (T - T_a) - \frac{4\sigma\varepsilon P T_a^3}{A_{cr}} (T - T_a) + (1-\phi) q_o''' [1 + \gamma(T - T_a)] = (\rho c_p)_{eff} u \frac{\partial T}{\partial x} + (\rho c_p)_{eff} \frac{\partial T}{\partial t} \tag{16}$$

when the like terms are collected, we have

$$\left(k_{eff} + \frac{16\sigma\phi T_a^3}{3\beta_R}\right) \frac{\partial^2 T}{\partial x^2} - \frac{\rho_f c_{p,f} g K \beta W \phi (T - T_a)^2}{v A_{cr}} - \left[ \frac{hP(1-\phi)}{A_{cr}} + \frac{4\sigma\varepsilon P T_a^3}{A_{cr}} \right] (T - T_a) - (\gamma(1-\phi) q_o''') (T - T_a) + (1-\phi) q_o''' = (\rho c_p)_{eff} u \frac{\partial T}{\partial x} + (\rho c_p)_{eff} \frac{\partial T}{\partial t} \tag{17}$$

Which can be written as

$$\left(1 + \frac{16\sigma\phi T_a^3}{3\beta_R k_{eff}}\right) \frac{\partial^2 T}{\partial x^2} - \frac{\rho_f c_{p,f} g K \beta W \phi (T - T_a)^2}{v A_{cr} k_{eff}} - \left[ \frac{hP(1-\phi)}{A_{cr} k_{eff}} + \frac{4\sigma\varepsilon P T_a^3}{A_{cr} k_{eff}} \right] (T - T_a) - \frac{(\gamma(1-\phi) q_o''')}{k_{eff}} (T - T_a) + \frac{(1-\phi) q_o'''}{k_{eff}} = \frac{(\rho c_p)_{eff} u}{k_{eff}} \frac{\partial T}{\partial x} + \frac{(\rho c_p)_{eff}}{k_{eff}} \frac{\partial T}{\partial t} \tag{18}$$

Adopting the adimensional variables below

$$X = \frac{x}{L}, \quad \theta = \frac{T - T_a}{T_b - T_a}, \quad \tau = \frac{k_{eff} t}{(\rho c_p)_{eff} L^2} \tag{19}$$

The controlling equation in Eq. (18) becomes

$$(1+4R)\frac{\partial^2\theta}{\partial X^2} - S_h\theta^2 - (Mc + Mr - Q_{\gamma e})\theta + Q_e = Pe_e \frac{\partial\theta}{\partial X} + \frac{\partial\theta}{\partial\tau} \quad (20)$$

where

$$Rd = \frac{16\sigma\phi T_a^3}{3\beta_R k_{eff}}, \quad Mc^2 = \frac{hP(1-\phi)L^2}{A_{cr}k_{eff}}, \quad Mr = \frac{4\sigma\varepsilon PT_a^3L^2}{A_{cr}k_{eff}}, \quad Q_\gamma = \frac{\gamma(1-\phi)q_o''L^2}{k_{eff}}, \quad (21)$$

$$Q = \frac{(1-\phi)q_o''L^2}{k_{eff}(T_b - T_a)}, \quad Pe = \frac{(\rho c_p)_{eff} u_o L}{k_{eff}}, \quad S_h = \frac{\rho_f c_{p,f} g K \beta W \phi L^2 (T_b - T_a)}{\nu A_{cr} k_{eff}}$$

Eq. (20) is alternative written as

$$\frac{\partial^2\theta}{\partial X^2} - S_p\theta^2 - (Nc + Nr - Q_\gamma)\theta + Q = Pe_R \frac{\partial\theta}{\partial X} + \zeta \frac{\partial\theta}{\partial\tau} \quad (22)$$

where

$$Nc = \frac{Mc}{1+4R}, \quad Nr = \frac{Mr}{1+4R}, \quad S_p = \frac{S_h}{1+4R}, \quad Q_\gamma = \frac{Q_{\gamma e}}{1+4R},$$

$$Q = \frac{Q_e}{(1+4R)}, \quad Pe_R = \frac{Pe_e}{(1+4R)}, \quad \zeta = \frac{1}{(1+4R)}$$

Therefore, the adimensional initial condition is

$$\theta = \theta_0 \text{ when } \tau = 0, \text{ for } 0 < X < 1, \quad (23)$$

and, the dimensionless boundary conditions are given as

$$\theta = 1, \text{ at } X = 0 \text{ for } \tau > 0, \quad (24a)$$

$$\frac{\partial\theta}{\partial X} = 0, \text{ at } X = 1, \text{ for } \tau > 0, \quad (24b)$$

### 3. METHOD OF LINES

Method of lines is a semi-discretization method. The inherent simplicity of the computational procedures and superiority of stability advantage over the direct finite difference method has encouraged its various applications. [37-74].

### 3. 1. Method of Lines the Transient Thermal Problems

Applying method of lines to Eq. (22)

$$\left(\frac{\theta_{i+1}-2\theta_i+\theta_{i-1}}{h^2}\right)-S_p\theta_i^2-(Nc+Nr-Q_\gamma)\theta_i+Q=Pe_R\left(\frac{\theta_{i+1}-\theta_{i-1}}{2h}\right)+\zeta\frac{d\theta_i}{d\tau} \quad (25)$$

Collecting like terms

$$\left(\frac{\theta_{i+1}-2\theta_i+\theta_{i-1}}{h^2}\right)-Pe_R\left(\frac{\theta_{i+1}-\theta_{i-1}}{2h}\right)-S_p\theta_i^2-(Nc+Nr-Q_\gamma)\theta_i+Q=\zeta\frac{d\theta_i}{d\tau} \quad (26)$$

Which can be expressed as

$$\frac{d\theta_i}{d\tau}=\left(\frac{\theta_{i+1}-2\theta_i+\theta_{i-1}}{h^2\zeta}\right)-Pe_R\left(\frac{\theta_{i+1}-\theta_{i-1}}{2h\zeta}\right)-\frac{S_p}{\zeta}\theta_i^2-\left(\frac{Nc+Nr-Q_\gamma}{\zeta}\right)\theta_i+\frac{Q}{\zeta} \quad (27)$$

$$i=1,2,\dots,N$$

From the initial condition is

$$\theta_i=0, \text{ at } \tau=0, \quad i=1,2,\dots,N \quad (28)$$

While for boundary conditions, we have

$$\frac{\theta_2(\tau)-\theta_0(\tau)}{2h}=0 \Rightarrow \theta_2(\tau)=\theta_0(\tau), \quad \rightarrow \theta_0(\tau)=\theta_0 \quad \text{at } i=0 \quad (29a)$$

and

$$\theta_{N+1}(\tau)=1, \quad \text{at } i=N+1 \quad (29b)$$

where

$N$  is the number of interior node points used in the discretization of the space,  $x \quad h = \frac{1}{N+1}$  is the node spacing or space step.

Which can be written as

$$\begin{aligned}
 \frac{d\theta_1}{d\tau} &= \left( \frac{\theta_2 - 2\theta_1 + \theta_0}{h^2\zeta} \right) - Pe_R \left( \frac{\theta_2 - \theta_0}{2h\zeta} \right) - \frac{S_p}{\zeta} \theta_1^2 - \left( \frac{Nc + Nr - Q_\gamma}{\zeta} \right) \theta_1 + \frac{Q}{\zeta} \\
 \frac{d\theta_2}{d\tau} &= \left( \frac{\theta_3 - 2\theta_2 + \theta_1}{h^2\zeta} \right) - Pe_R \left( \frac{\theta_3 - \theta_1}{2h\zeta} \right) - \frac{S_p}{\zeta} \theta_2^2 - \left( \frac{Nc + Nr - Q_\gamma}{\zeta} \right) \theta_2 + \frac{Q}{\zeta} \\
 \frac{d\theta_3}{d\tau} &= \left( \frac{\theta_4 - 2\theta_3 + \theta_2}{h^2\zeta} \right) - Pe_R \left( \frac{\theta_3 - \theta_1}{2h\zeta} \right) - \frac{S_p}{\zeta} \theta_3^2 - \left( \frac{Nc + Nr - Q_\gamma}{\zeta} \right) \theta_3 + \frac{Q}{\zeta} \\
 \frac{d\theta_N}{d\tau} &= \left( \frac{\theta_{N+1} - 2\theta_N + \theta_{N-1}}{h^2\zeta} \right) - Pe_R \left( \frac{\theta_{N+1} - \theta_{N-1}}{2h\zeta} \right) - \frac{S_p}{\zeta} \theta_N^2 - \left( \frac{Nc + Nr - Q_\gamma}{\zeta} \right) \theta_N + \frac{Q}{\zeta}
 \end{aligned} \tag{30}$$

The resulting ODE is generally of the form

$$\frac{d\theta_i}{dt} = A\theta_{i+1}^2 + B\theta_{i+1} + C\theta_i + D\theta_{i-1} + E \tag{31}$$

Eq. (18) is a system of  $N$  nonlinear first order differential equations and can be written in matrix form as

$$\frac{d\Theta}{dt} = A_M \Phi + A_L \Theta + K \tag{32}$$

where

$$A_M = A = \frac{S_p}{\zeta} \begin{bmatrix} -1 & 0 & 0 & \cdots & 0 & 0 & 0 \\ 0 & -1 & 0 & \cdots & 0 & 0 & 0 \\ 0 & 0 & -1 & \cdots & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & \cdots & -1 & 0 & 0 \\ 0 & 0 & 0 & \cdots & 0 & -1 & 0 \\ 0 & 0 & 0 & \cdots & 0 & 0 & -1 \end{bmatrix}$$



$$A_L = (B + C + D) = \frac{1}{h^2 \zeta} \begin{bmatrix} -2 & 1 & 0 & \dots & 0 & 0 & 0 \\ 1 & -2 & 1 & \dots & 0 & 0 & 0 \\ 0 & 1 & -2 & \dots & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & \dots & -2 & 1 & 0 \\ 0 & 0 & 0 & \dots & 1 & -2 & 1 \\ 0 & 0 & 0 & \dots & 0 & 1 & -2 \end{bmatrix} + \frac{Pe_R}{2h\zeta} \begin{bmatrix} 0 & 1 & 0 & \dots & 0 & 0 & 0 \\ 1 & 0 & 1 & \dots & 0 & 0 & 0 \\ 0 & 1 & 0 & \dots & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & \dots & 0 & 1 & 0 \\ 0 & 0 & 0 & \dots & 1 & 0 & 1 \\ 0 & 0 & 0 & \dots & 0 & 1 & 0 \end{bmatrix}$$

$$+ \left( \frac{Nc + Nr - Q_\gamma}{\zeta} \right) \begin{bmatrix} -1 & 0 & 0 & \dots & 0 & 0 & 0 \\ 0 & -1 & 0 & \dots & 0 & 0 & 0 \\ 0 & 0 & -1 & \dots & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & \dots & -1 & 0 & 0 \\ 0 & 0 & 0 & \dots & 0 & -1 & 0 \\ 0 & 0 & 0 & \dots & 0 & 0 & -1 \end{bmatrix}$$

$$\bar{\Phi} = [\theta_1^2 \quad \theta_2^2 \quad \dots \quad \theta_N^2]^T$$

$$\Theta = [\theta_1 \quad \theta_2 \quad \dots \quad \theta_N]^T$$

and

$b$  is a column vector of order  $N \times 1$  which is in the form

$$E = \frac{Q}{\zeta} [1 \quad 1 \quad \dots \quad 1 \quad 1]^T$$

Euler's method is used to solve the system of the nonlinear differential equations. The method is stated as follows

$$\theta_i^{n+1} = \theta_i^n + hf(\theta_i^n)$$

where  $f(\theta_i^n) = A\theta_{i+1}^2 + B\theta_{i+1} + C\theta_i + D\theta_{i-1} + E$

$h$  is the time-step

For the linear thermal model, the resulting ODE is generally of the form

$$\frac{d\theta_i}{dt} = B\theta_{i+1} + C\theta_i + D\theta_{i-1} + E \tag{33}$$

Also, Eq. (33) is a system of  $N$  linear first order differential equations and can be written in matrix form as

$$\frac{d\Theta}{dt} = A_L \Theta + b \tag{34}$$

where,

$A_L$  is a  $N \times N$  matrix given by

$$A_L = \frac{1}{h^2} \begin{bmatrix} -2 & 1 & 0 & \dots & 0 & 0 & 0 \\ 1 & -2 & 1 & \dots & 0 & 0 & 0 \\ 0 & 1 & -2 & \dots & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & \dots & -2 & 1 & 0 \\ 0 & 0 & 0 & \dots & 1 & -2 & 1 \\ 0 & 0 & 0 & \dots & 0 & 1 & -2 \end{bmatrix}$$

$$(B+C+D) = \frac{1}{h^2 \zeta} \begin{bmatrix} -2 & 1 & 0 & \dots & 0 & 0 & 0 \\ 1 & -2 & 1 & \dots & 0 & 0 & 0 \\ 0 & 1 & -2 & \dots & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & \dots & -2 & 1 & 0 \\ 0 & 0 & 0 & \dots & 1 & -2 & 1 \\ 0 & 0 & 0 & \dots & 0 & 1 & -2 \end{bmatrix} + \frac{Pe_R}{2h\zeta} \begin{bmatrix} 0 & 1 & 0 & \dots & 0 & 0 & 0 \\ 1 & 0 & 1 & \dots & 0 & 0 & 0 \\ 0 & 1 & 0 & \dots & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & \dots & 0 & 1 & 0 \\ 0 & 0 & 0 & \dots & 1 & 0 & 1 \\ 0 & 0 & 0 & \dots & 0 & 1 & 0 \end{bmatrix}$$

$$+ \left( \frac{Nc + Nr - Q_\gamma}{\zeta} \right) \begin{bmatrix} -1 & 0 & 0 & \dots & 0 & 0 & 0 \\ 0 & -1 & 0 & \dots & 0 & 0 & 0 \\ 0 & 0 & -1 & \dots & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & \dots & -1 & 0 & 0 \\ 0 & 0 & 0 & \dots & 0 & -1 & 0 \\ 0 & 0 & 0 & \dots & 0 & 0 & -1 \end{bmatrix}$$

and

$$\Theta = [\theta_1 \quad \theta_2 \quad \dots \quad \theta_N]^T$$

$b$  is a column vector of order  $N \times 1$  which is in the form

$$b = \frac{1}{h^2} [\theta_0 \quad 0 \quad \dots \quad 0 \quad \theta_{N+1}]^T = \frac{1}{h^2} [\theta_0 \quad 0 \quad \dots \quad 0 \quad 1]^T$$

Eigen value and Eigen vector method is used to solve the system of the linear differential equations

### 3. NUMERICAL SOLUTION OF THE NONLINEAR TRANSIENT THERMAL MODEL

$$\begin{aligned} & \left( \frac{\theta_{i+1}^{n+1} - 2\theta_i^{n+1} + \theta_{i-1}^{n+1} + \theta_{i+1}^n - 2\theta_i^n + \theta_{i-1}^n}{2\Delta^2 X} \right) - S_p (\theta_i^n)^2 - (Nc + Nr - Q_\gamma) \theta_i^n + Q \\ & = Pe_R \left( \frac{\theta_{i+1}^{n+1} - \theta_{i-1}^{n+1} + \theta_{i+1}^n - \theta_{i-1}^n}{4\Delta X} \right) + \zeta \left( \frac{\theta_i^{n+1} - \theta_i^n}{\Delta \tau} \right) \end{aligned} \tag{35}$$

Which can be arranged as

$$\begin{aligned} & \left( \frac{1}{2\Delta^2 X} - \frac{Pe_R}{4\Delta X} \right) \theta_{i+1}^{n+1} - \left( \frac{2}{2\Delta^2 X} + \frac{\zeta}{\Delta \tau} \right) \theta_i^{n+1} + \left( \frac{1}{2\Delta^2 X} + \frac{Pe_R}{4\Delta X} \right) \theta_{i-1}^{n+1} + \left( \frac{1}{2\Delta^2 X} - \frac{Pe_R}{4\Delta X} \right) \theta_{i+1}^n \\ & - \left( \frac{2}{2\Delta^2 X} - \frac{\zeta}{\Delta \tau} + (Nc + Nr - Q_\gamma) \right) \theta_i^n - S_p (\theta_i^n)^2 + \left( \frac{1}{2\Delta^2 X} + \frac{Pe_R}{4\Delta X} \right) \theta_{i-1}^n + Q = 0 \end{aligned} \tag{36}$$

We can write Eq. (36) as

$$A\theta_{i+1}^{n+1} + B\theta_i^{n+1} + C\theta_{i-1}^{n+1} + D\theta_{i+1}^n + E\theta_i^n + F(\theta_i^n)^2 + G\theta_{i-1}^n + H = 0 \tag{37}$$

where

$$\begin{aligned} A &= \left( \frac{1}{2\Delta^2 X} - \frac{Pe_R}{4\Delta X} \right), \quad B = -\left( \frac{2}{2\Delta^2 X} + \frac{\zeta}{\Delta \tau} \right), \quad C = \left( \frac{1}{2\Delta^2 X} + \frac{Pe_R}{4\Delta X} \right), \quad D = \left( \frac{1}{2\Delta^2 X} - \frac{Pe_R}{4\Delta X} \right), \\ E &= -\left( \frac{2}{2\Delta^2 X} - \frac{\zeta}{\Delta \tau} + (Nc + Nr - Q_\gamma) \right), \quad F = -S_p, \quad G = \left( \frac{1}{2\Delta^2 X} + \frac{Pe_R}{4\Delta X} \right), \quad H = Q, \end{aligned}$$

The finite difference discretization of the initial condition is

$$\theta_i^o = 0, \tag{38}$$

While FDM for boundary conditions become

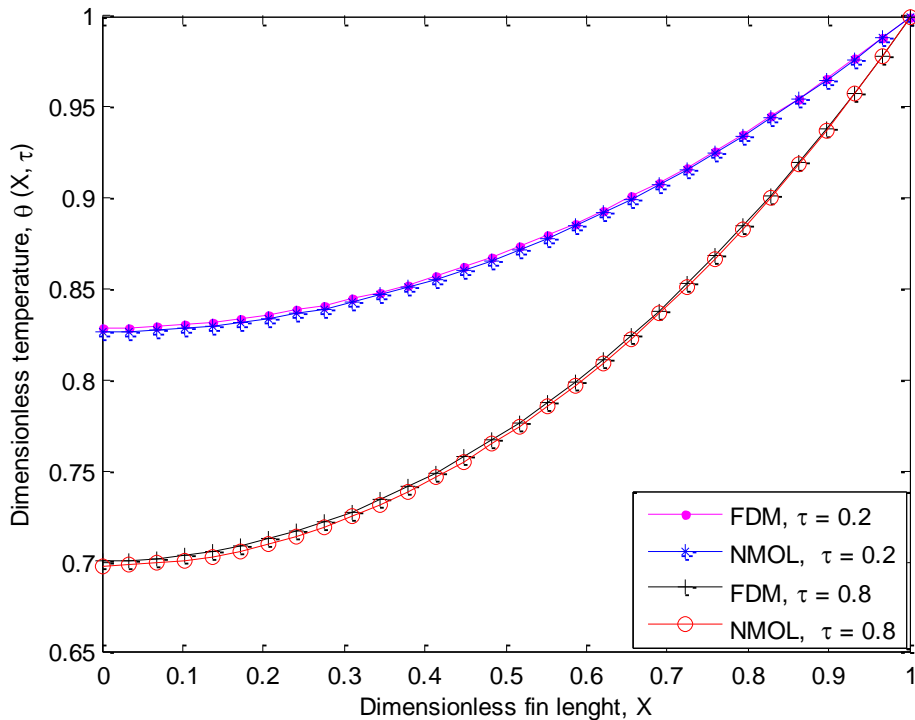
$$\frac{\theta_1^n - \theta_{-1}^n}{2\Delta X} = 0 \Rightarrow \theta_1^n = \theta_{-1}^n \tag{39}$$

$$\theta_M^n = 1$$

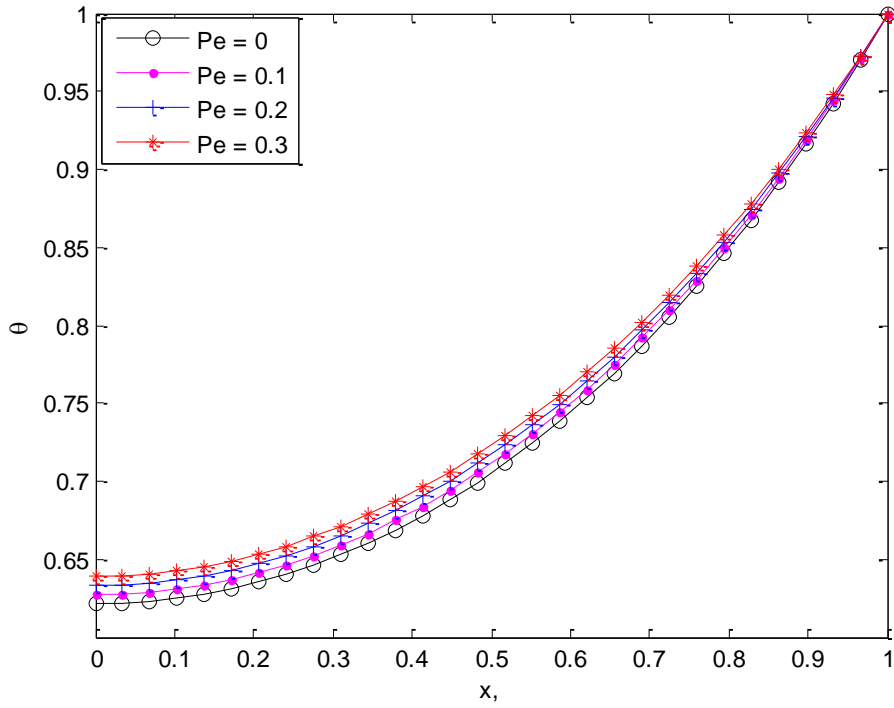
#### 4. RESULTS AND DISCUSSION

In order to establish the accuracy of the present method, we compared the results of the numerical method of lines (NMOL) with the results of finite difference method (FDM) as shown in Fig. 2. Having established this fact, the numerical solutions was then used to explored the impact of Peclet number on the thermal responses of the extended surfaces under various values of the controlling parameters.

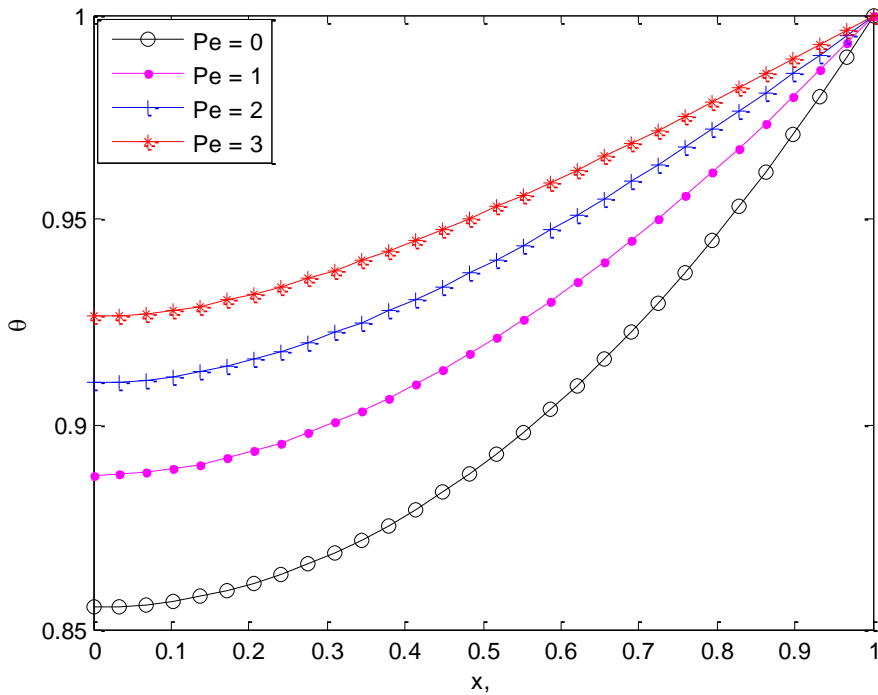
Fig. 3 presents the effect of Peclet number on the thermal response of the fin when it is not moving and it is moving slowly. The Peclet number bears direct relationship with the fin temperature as depicted in the figure. This is due to the increase in the speed of the fin that reduces the material exposure time to the surrounding fluid. Therefore, low Peclet number favours cooling enhancement.



**Fig. 2.** Comparison of results of NMOL and FDM



**Fig. 3.** Effect of Peclet number on the temperature when  $Nc=0.8$ ,  $Nr=0.2$ ,  $\theta_a=0.1$   $\theta_s=0.2$



**Fig. 4.** Impact of Peclet number on the temperature when  $Nc=0.5$ ,  $Nr=0.5$   $\theta_a=0.3$   $\theta_s=0.2$

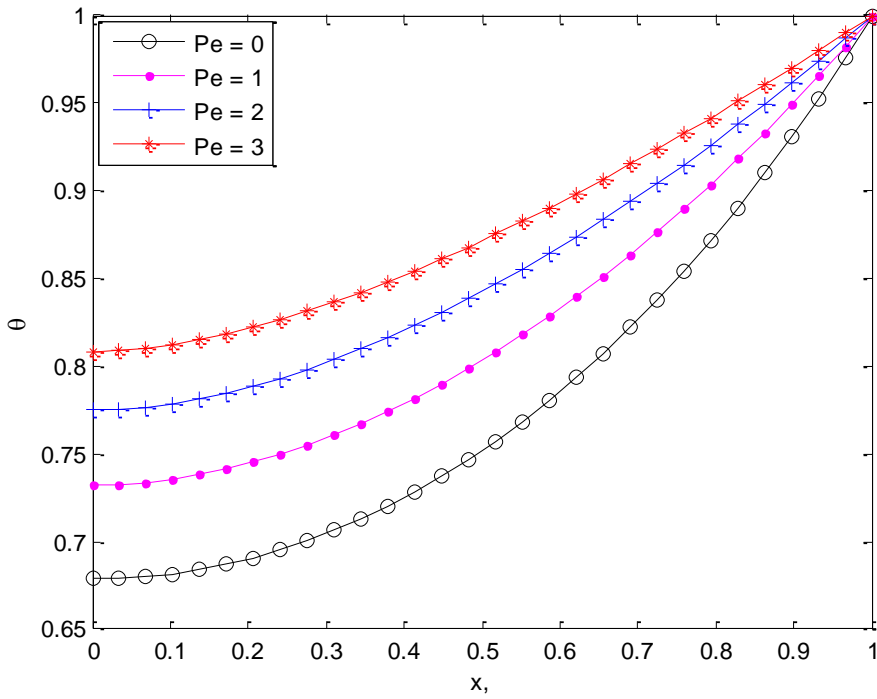


Fig. 5. Impact of Peclet number on the temperature when  $Nc=0.5$ ,  $Nr=0.8$   $\theta_a=0.2$ ,  $\theta_s=0.2$

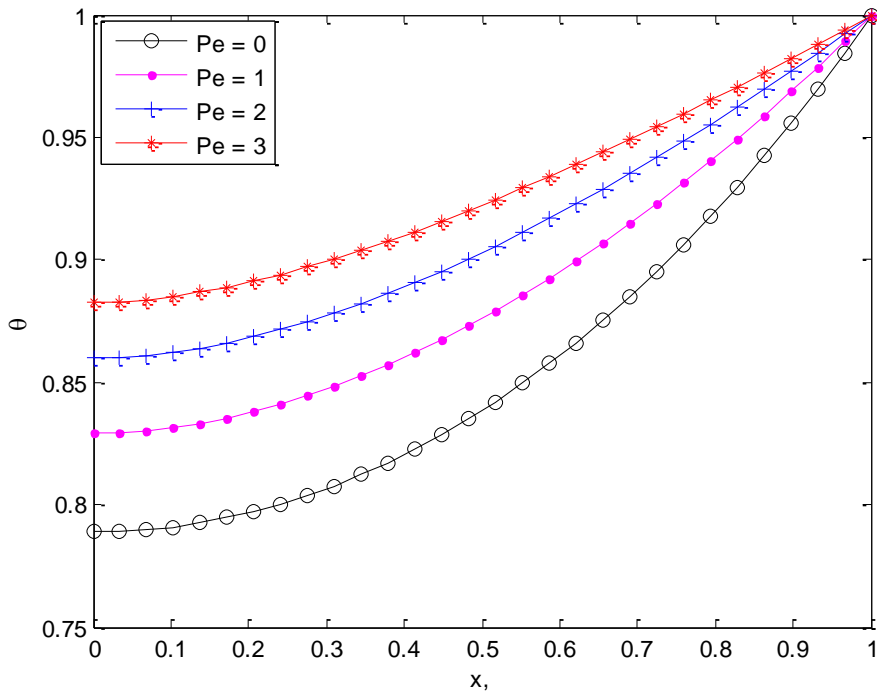
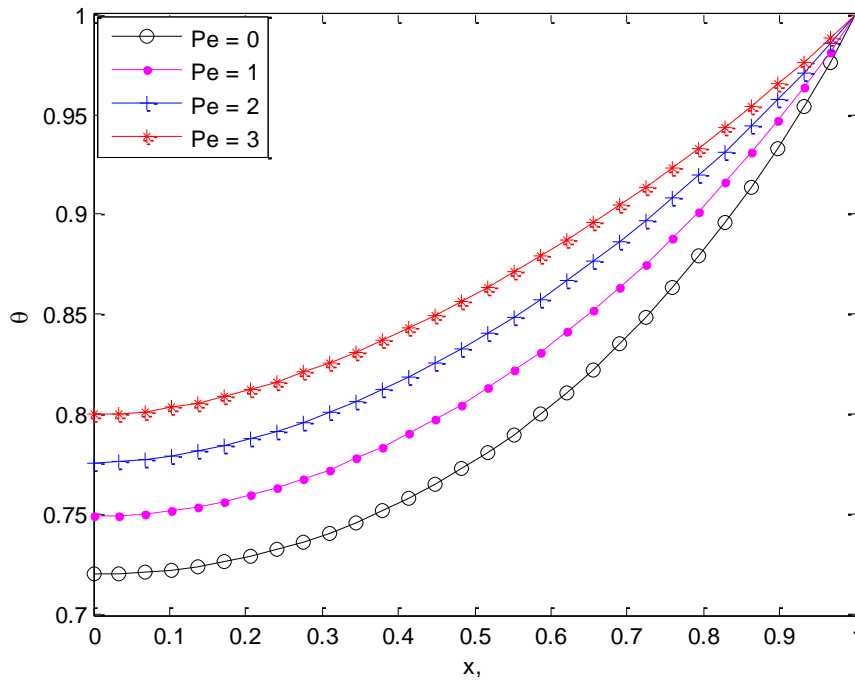
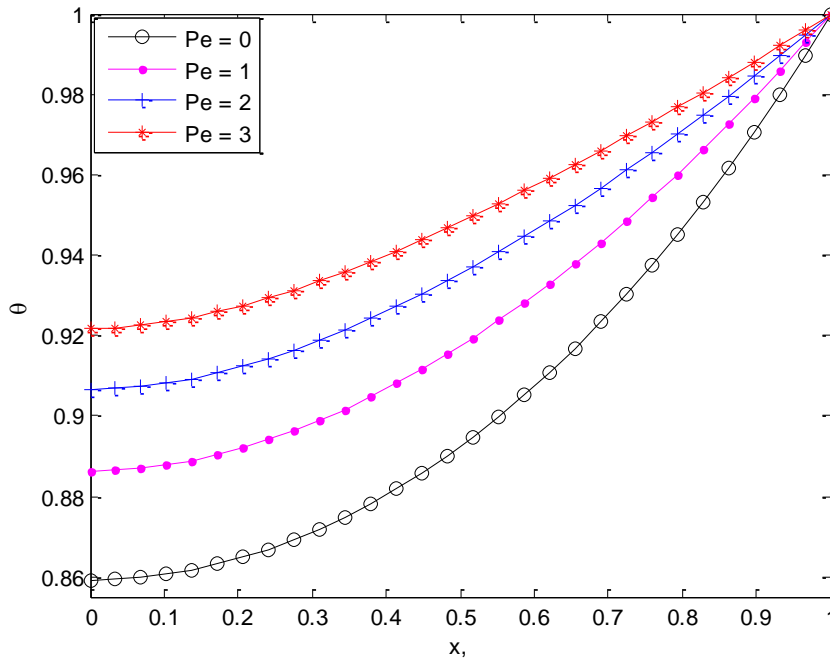


Fig. 6. Effect of Peclet number on the temperature when  $Nc=0.5$ ,  $Nr=0.5$   $\theta_a=0.2$   $\theta_s=0.4$



**Fig. 7.** Impact of Peclet number on the temperature when  $Nc=0.5$ ,  $Nr=0.5$ ,  $S_p=0.3$ ,  $\theta_a=0.2$ ,  $\theta_s=0.0$



**Fig. 8.** Impact of Peclet number on the temperature when  $Nc=0.8$ ,  $Nr=0.2$ ,  $Q = 0.1$ ,  $\theta_a=0.2$ ,  $\theta_s=0.2$

Figs. 4-8 illustrate the impact of Peclet number on the thermal responses of the extended surfaces under various values of the controlling parameters. The controlling parameters such as convective-conductive, radiative-conductive, ambient temperature, porosity and internal heat generation parameters have significant influence on the thermal behaviour of the porous moving fin. Closer investigation revealed that when the convective-conductive and radiative-conductive parameters increase, the fin temperature decreases. However, the fin Peclet number increases when ambient and surface temperatures increase while the temperature of the fin decreases as the convective-conductive, radiative-conductive and porous terms increase. Additionally, an increase in the internal heat generation causes the fin temperature to increase. Under all these various varying values of the controlling parameters, it is established that the temperature of the fin increases as the Peclet number increases as depicted in all the figures.

## **5. CONCLUSION**

Numerical method of lines has been applied in this study to investigate the effect of Peclet number on the transient nonlinear thermal behaviour of moving porous fins. The results showed that the Peclet number increases with increasing fin temperature. The numerical investigations reveal that the fin temperature decreases as the convective-conductive, radiative-conductive and porous terms increase. However, when internal heat generation, ambient and surface temperatures increase, the fin Peclet number also increases. Under all these various varying values of the controlling parameters the temperature of the fin increases as the Peclet number increases. The physical insight depicted in this study will create a new path for the nonlinear analysis of the thermal problems.

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